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Scanning Tunneling Microscopy and Medium Energy Ion Scattering Spectrometry of Spinel Structure of $\text{Li}_4\text{Ti}_5\text{O}_{12}$ Surface

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Spinel lithium titanate ($\text{Li}_4\text{Ti}_5\text{O}_{12}$) is one of the promising anode materials for high-performance lithium-ion batteries (LIBs). It is crucial to investigate atomistic structures of $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surfaces to understand the phenomena at the LTO/electrolyte interfaces such as CO_2 -gas generation which greatly affects the performance and safety of LIBs. By applying scanning tunneling microscopy (STM) and medium energy ion scattering (MEIS) to a $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) film prepared from a TiO_2 wafer, we found that there exists two kinds of Li-terminated (111) terraces bounded by steps with different heights. In the major terraces, the top hexagonal Li layer is stacked above the oxygen layer, while the top Li layer is stacked above the Ti-Li layer in the minor terraces. The relative stability between the two surface structures seems to depend on the atmosphere due to different stoichiometry, which is a possible origin of CO_2 generation via redox interaction with electrolyte molecules.

Keywords: $\text{Li}_4\text{Ti}_5\text{O}_{12}$; Spinel structures; Scanning tunneling microscopy; MEIS
Received: 15 May 2018, **Revised:** 18 July 2018; **Accepted:** 25 July 2018

1. Introduction

Lithium transition-metal oxides [1-3] are widely used as active electrode materials for lithium-ion batteries (LIBs) due to their reversible lithium insertion/extraction properties. Especially, the spinel lithium titanium oxide ($\text{Li}_4\text{Ti}_5\text{O}_{12}$, LTO) [3] is one of the most important anode materials for its high rate Li insertion/extraction (charge-discharge) capability [4,5] with excellent durability[6]. These properties, based on the rigid lattice stability of the spinel framework [3-5], enable the design of a safer anode of LIBs than a carbon-based one, promising for the usage in electric vehicles and other wide industrial applications [7]. Numerous studies have been performed on the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ bulk structure and its changes by Li insertion/extraction [8-13], while the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surface structure has not yet been understood. The Li insertion/extraction occurs via the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ /electrolyte interface, and the interface phenomena such as CO_2 -gas generation [14,15] of $\text{Li}_4\text{Ti}_5\text{O}_{12}$ in high-power LIBs. Therefore, the investigation of atomistic $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surface structure is essential to understand the atomistic mechanisms of electrochemical processes and various phenomena at $\text{Li}_4\text{Ti}_5\text{O}_{12}$ /electrolyte interfaces [16-18], seriously affecting the performance and safety of LIBs. Except for graphite electrodes, there are few reports on atomic scale structures of surfaces of electrode materials, because of the difficulty in preparing crystalline samples with atomically-flat surfaces suitable for scanning-probe microscopy (SPM). Recently, we developed the technique to prepare $\text{Li}_4\text{Ti}_5\text{O}_{12}$ films with atomically-flat surfaces on a TiO_2 wafer using conventional solid-state reaction

[19]. The prepared $\text{Li}_4\text{Ti}_5\text{O}_{12}$ film was shown to have electrochemical activity essentially similar to usual $\text{Li}_4\text{Ti}_5\text{O}_{12}$ power. For the film sample, the morphology changes of a $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface in the early stage of initial Li insertion process [20] were investigated by using atomic force microscopy (AFM).

In the present study, we investigated the atomic scale structure of a $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface by using high-resolution scanning tunneling microscopy (STM). We also apply medium energy ion scattering (MEIS) analysis, which can provide complementary information on the elemental depth profile in the surface and sub-surface layers. Atomic models from the experimental results are discussed from the viewpoints of surface stoichiometry and surface electronic structure.

2. Experimental Part

A $\text{Li}_4\text{Ti}_5\text{O}_{12}$ thin-film sample was prepared by solid state reaction as previously reported [19]. A commercial rutile- TiO_2 (111) wafer (SHINKOSHA, Co. LTD) was cleaned in analytical grade acetone by an ultrasonic cleaner to remove surface contaminations. The wafer with a size of $2 \times 6 \times 0.5$ mm was sealed in a 99% alumina crucible with 1 mg of $\text{LiOH} \cdot \text{H}_2\text{O}$ and calcined for 15 h at 1173 K in air. In the calcination process, LiOH is vaporized, and a single-crystalline $\text{Li}_4\text{Ti}_5\text{O}_{12}$ film with about $15 \mu\text{m}$ thickness is grown on the TiO_2 (111) surface. The crystal structure of the grown film was evaluated by a Mo-K α X-ray diffractometer (X'Pert PRO, PANalytical) with out-of-plane measurement. AFM (NanoNavi-II, SIL) was performed by dynamic force

mode in air at 1000nm^2 area. We observed the surface structure in an atomic scale with high-resolution STM (JSTM-4500TX, JEOL) under ultra-high vacuum (UHV, $\sim 10^7$ Pa) at room temperature. Before the STM observation, the sample was heated by silicone heat board at 1073 K for 1 h and 1123 K for 15 min under the UHV to prepare the clean surface. STM images were acquired at +1 to +1.5 V of sample bias and 0.3nA of tunneling current. High-resolution MEIS using 80 keV He^+ ions allowed for elemental depth profiling, which was carried out at beam-line 8 named SORIS at Ritsumeikan SR Center. Prior to the MEIS analysis, we heated the sample at 923 K for 1 h and then at 1123 K for 15 min in UHV and confirmed no carbon contamination by Auger electron spectroscopy (see Fig. 1).

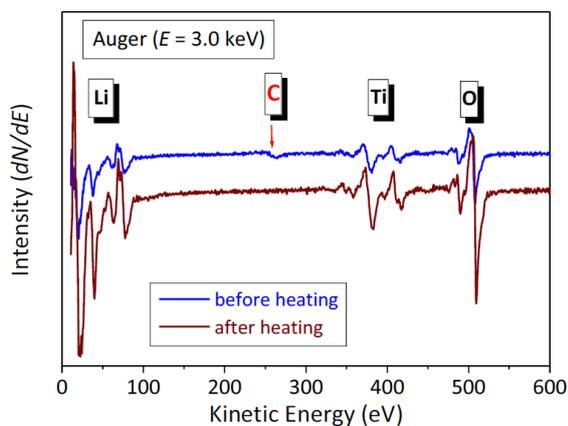


Fig. (1) Auger spectra taken for the sample before (blue) and after heating in UHV

3. Results and Discussion

Here, we explain the (111) atomic-layer stacking in a spinel $\text{Li}_4\text{Ti}_5\text{O}_{12}$ crystal ($Fd-3m$, $a = 0.8357$ nm [21]) before reporting experimental results. Figure 2(a) shows four kinds of (111) atomic layers. Li atoms at the 8a sites of Wyckoff positions form the 8a (Li) layer (indicated by a red solid line), and O atoms at the 32e sites form the 32e (O) layer (indicated by a black solid line). The 16d-1 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) and 16d-2 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layers (indicated by blue and green solid lines) consist of Li and Ti atoms with the ratio of 1:5 at the 16d sites. As shown in Fig. (2b), the in-plane atomic density is equal to that of the 8a (Li) and 16d-1 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layers, while it is three times larger that of the 16d-2 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layer and four times larger that of the 32e (O) layer. The 8a (Li) and 16d-1 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layers should have hexagonal atomic arrangements with the interatomic distance of about 0.6 nm, while the minimum interatomic distance in the 16d-2 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layer is about 0.3 nm. Note that the distinction between Ti and Li atoms on the 16d-1 and 16d-2 layers is not definite essentially, while only the ratio of 1(Li):5(Ti) is definite. In the (111) stacking sequence in Fig. (2a), total six layers as two 8a (Li), two 32e (O), one 16d-1 and one 16d-2 layers, constitute one period.

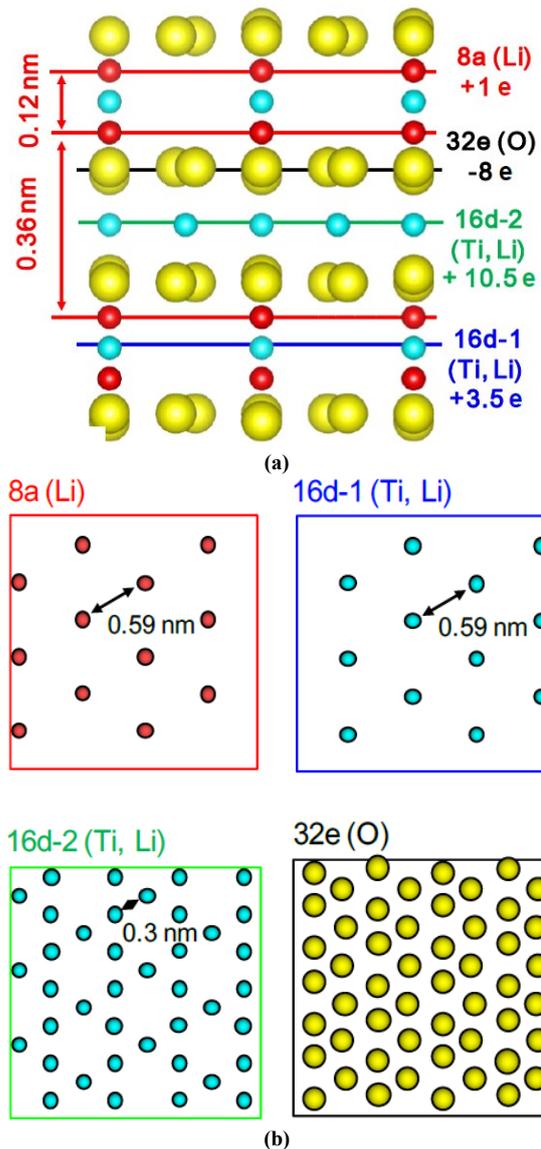


Fig. (2) (a) Out-of-plane structure of (111) layer stacking in $\text{Li}_4\text{Ti}_5\text{O}_{12}$ viewed from [112] direction. Formal ionic charges of each atomic layer per 1×1 period are denoted. (b) In-plane atomic arrangements of four kinds of (111) layers shown in (a). Values of intervals between the 8a (Li) layers and interatomic distances in the hexagonal atomic arrangements are estimated from the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ crystal structure with the experimental lattice constant of $a = 0.8357$ nm [21]

From the crystallographic data of the bulk $\text{Li}_4\text{Ti}_5\text{O}_{12}$ [21], it can be said that only the 8a (Li) layers are repeated with the two kinds of intervals as about 0.12 and 0.36 nm. In Fig. (2a), a formal ionic charge of each (111) atomic layer per 1×1 hexagonal period is also denoted. This is obtained by the sum of charges of each species as $+1e + 4e$ (Ti) and $2e$ (O) on each layer. It can be said that each oxygen layer with $8e$ (four oxygen atoms per 1×1 period) receives $5.25e$ from the lower 16d-2 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layer and $2.75e$ from the upper 8a (Li) and 16d-1 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layers.

In Fig. 3(a), the out-of-plane X-ray (Mo- $K\alpha$) diffraction spectrum of the prepared film reveals strong and sharp diffraction lines of the $\text{Li}_4\text{Ti}_5\text{O}_{12}$

(111) series, indicating that the prepared film is highly-ordered (111) crystalline film. Fig. 3(b) shows the AFM image of the prepared (111) film surface. We can see clear steps with a single atomic-layer height of 0.48 nm and atomically-flat terraces with $Ra = 0.02$ nm of surface roughness. This step height corresponds to one period of the sequence of (111) atomic layers in Fig. 2(a).

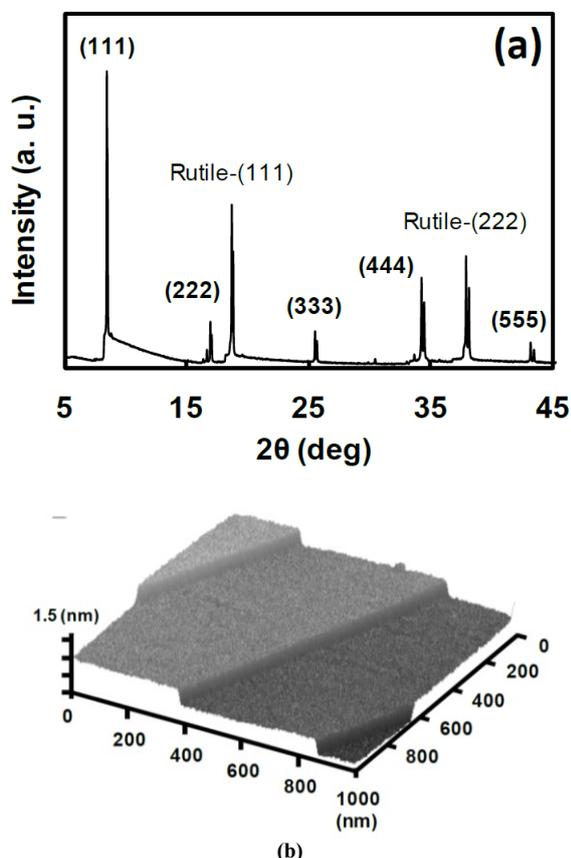


Fig. (3) (a) Out-of-plane X-ray diffraction spectrum of the prepared $\text{Li}_4\text{Ti}_5\text{O}_{12}$ film. (b) AFM image of the prepared $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface acquired from 1000×1000 nm scan area

Figure 4(a) shows the STM image of the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface acquired from 50×50 nm area, where we can see clear steps and terraces with some atomistic defects. This image is not the periodic one of some artifact, but is the true image of atomic scale resolution. Figure 4(b) shows the surface height profile of X-Y line across the two steps in Fig. 4(a). The heights of the two steps are 0.35 and 0.13 nm. This indicates the presence of two kinds of different (111) terraces, in contrast to the AFM results of the single step height of 0.48 nm in Fig. 3(b), indicating only one type of (111) terrace. Figure 4(c) shows the high-resolution STM image acquired from the square area in the upper terrace in Fig. 4(a). We can see a hexagonal arrangement of bright spots with intervals of about 0.6 nm, while there also seem to exist some defects. A similar hexagonal arrangement of spots with similar intervals is also observed on the middle and lower

terraces in Fig. 4(a).

The observed hexagonal atomic arrangement of bright spots on each terrace indicates that the top layers of the two kinds of terraces should be $8a$ (Li) or $16d-1$ ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layers shown in Fig. (2b), having a similar hexagonal arrangement with interatomic distance of about 0.6 nm.

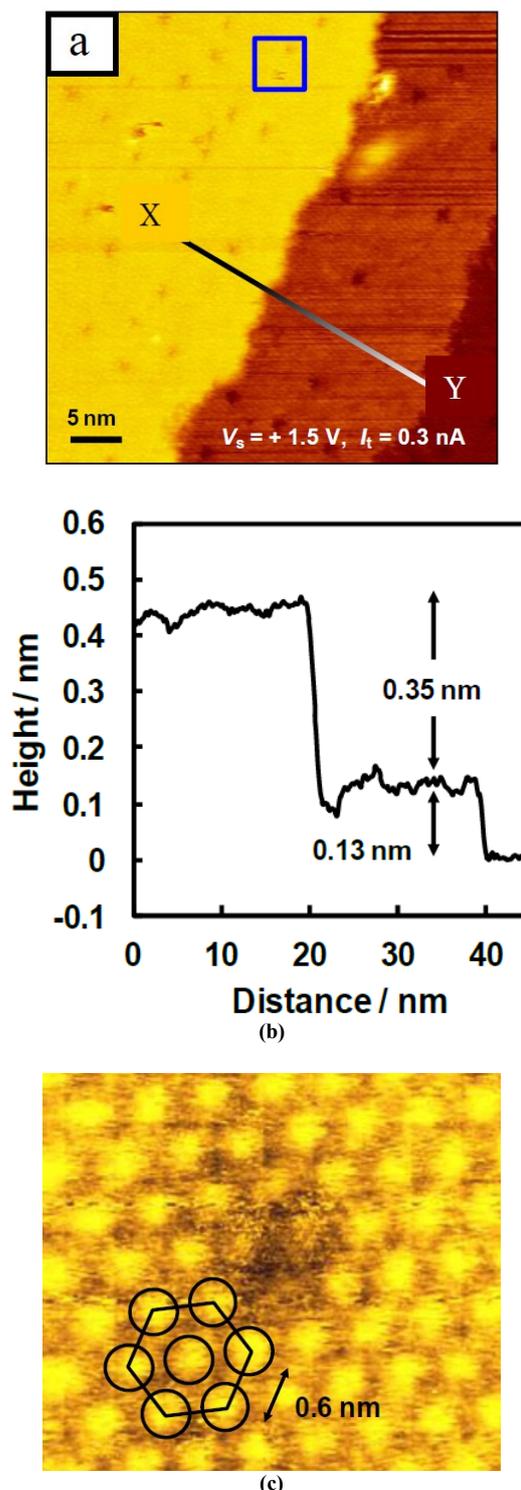


Fig. (4) (a) Atomic scale image of the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface acquired by UHV-STM (JEOL-4500TX). (b) Surface height profile of X-Y line shown in (a). (c) Magnified STM image of the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface acquired from the square area in (a)

From the step heights, we can expect that both kinds of terraces are Li-terminated surfaces as the type (A) and type (B) surfaces shown in Fig. (5). The upper terrace in Fig. (4a) should consist of the 8a (Li) layer stacked on the 32e (O) layer (type (B)), while the middle terrace should consist of the 8a (Li) layer stacked on the 16d-1 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layer (type (A)). Only this combination can satisfy the relation between the observed step heights (0.13 nm and 0.35 nm) and the intervals between the atomic layers in Fig. (2a) within the models of the (111) atomic-layer sequence similar to the bulk. It is apparent that the two models of the Li-terminated surfaces in Fig. 5 have different chemical compositions (stoichiometry) in the surface regions. As analyzed later, the type (B) surface is oxygen-rich and the type (A) surface is oxygen-deficient. Thus it is natural that the two kinds of the Li-terminated surfaces have different stabilities depending on the atmosphere, namely the oxygen chemical potential, as analyzed by Gibbs free energies.

In order to prove the present atomic models of the two kinds of (111) terraces, we have performed the high-resolution MEIS observation. This technique can identify a target atomic mass M with a resolution of $\Delta M/M \sim 1/100$ and provide valuable information on the atomic species and their arrangement in the surface and sub-surface layers. The details of MEIS analysis were described elsewhere[22,23]. The sample was prepared in the same manner as described for the STM observation and thus annealed at 923 K for 1 h and then at 1123 K for 15 min in UHV, resulting in complete elimination of carbon contaminations, as indicated in Fig. 1. Figure 6 shows the MEIS spectrum (open circles) observed for 80 keV He^+ ions incident on the clean $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface and scattered from the Ti atoms. The vertical arrow in this figure indicates the energy position of He^+ ions scattered from Ti atoms on top of the surface assumed to exist. The slightly lower energy shift of the leading edge (mid energy position of the slope) suggests the presence of a few over-layers consisting of low Z-number atoms above the Ti atoms. In other words, there are no Ti atoms on top of any terraces of the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface. The blue and red solid lines in Fig. 6 show the simulated spectra assuming the type (A) and type (B) surfaces, respectively. The small surface peak appearing at 67.5 keV in the blue line comes from the 16d-1 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layer below the 8a (Li) layer in the type (A) surface. In contrast, for the type (B) surface, the presence of two over-layers, the 8a (Li) and 32e (O) layers, smears out such a surface peak from the 16d-2 ($\text{Li}_{1/6}\text{Ti}_{5/6}$) layer, because of increased energy loss and straggling for He^+ ions subjected during passing through the over-layers. Note that the energy loss and straggling are caused by the interactions between He^+ ions and target electrons and thus lower Z-number materials

give lower energy loss and straggling. It is clear that the type (B) surface structure gives the best-fit to the observed MEIS spectrum. However, the MEIS results do not necessarily rule out the possibility of the coexistence of the type (A) surface as the minority portion of terraces. In Fig. (6), the green line shows the simulated spectrum for the 1:1 mixture of the type (A) and type (B) surfaces, where the effect of the type (A) surface is observed as a small shoulder. Also in the experimental spectrum, there seems to exist a very small shoulder, while it might be caused by some kind of errors or fluctuations. We have also performed simulations for smaller portions of the type (A) surface less than 50%, which indicates no clear effects in the simulated spectra from the coexistence of the type (A) surface less than $\sim 20\%$. Thus the MEIS results cannot exclude the possibility of the coexistence of the type (A) surface less than 20% for the sample analyzed by MEIS.

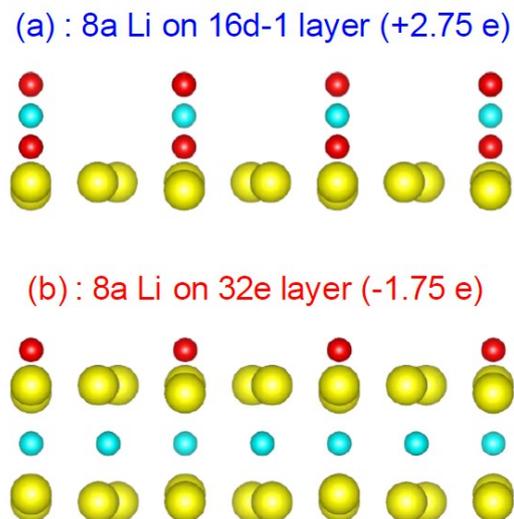


Fig. (5) Side views of two types of Li-terminated $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surfaces seen from the [112] direction. Red, blue and yellow circles denote, respectively lithium, titanium and oxygen atoms

The present MEIS result indicating that the type (B) surface is a large majority in the sample surface is consistent with the fact that the area of the type (B) surface as the upper terrace is much wider than the area of the type (A) surface as the middle terrace in the STM image of Fig. (4a), and the fact that occasionally only one type of the (111) terrace, probably the type (B) structure, with the single step height of 0.48 nm is observed in the STM image. Note that the high-temperature treatment in UHV for the MEIS sample is similar to the sample observed by STM.

The present models of two kinds of Li-terminated surfaces have the sequence of (111) atomic-layer stacking similar to the bulk. Of course, we do not deny general possibilities of further reconstruction in in-plane or stacking structure at the

surface. The present experimental results of step heights, periodic bright spots in STM images and MEIS naturally leads to the present models and we did not obtain any results indicating further reconstruction. About cation termination in stable oxide surfaces in usual atmosphere, note that there exist a lot of examples similar to the present type (B) structure, such as spinel $\text{Co}_3\text{O}_4(111)$ and $\text{Fe}_3\text{O}_4(111)$ surfaces [24-28] and corundum $\text{Al}_2\text{O}_3(0001)$ surface [29], where the cation topmost layer is stacked above the surface oxygen layer with the atomic-layer stacking sequence similar to the bulk. A recent study of the $\text{Li}_2\text{TiO}_3(001)$ surface with components similar to $\text{Li}_4\text{Ti}_5\text{O}_{12}$ has also concluded the cation (Li^+) termination from the results of STM observation and classical molecular dynamics simulations[30]. Similarly, we think that unoccupied orbitals at surface cations (Li^+) on the type (B) surface should be concerned with bright spots in the STM image.

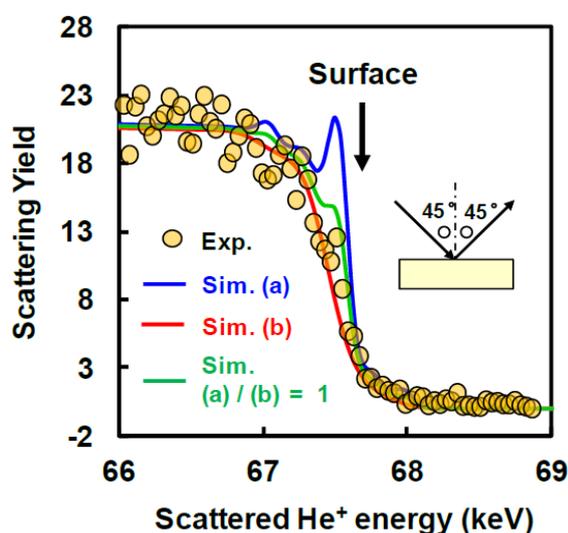


Fig. (6) MEIS spectrum observed for 80keV He^+ ions incident on the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface at 45° with respect to surface normal and scattered to 45° from Ti atoms. Blue, red and green curves are simulated MEIS spectra assuming type (A), type (B) and their 1:1 mixed surfaces, respectively

As mentioned above, the type (A) and type (B) surface structures have different surface stoichiometries. We can define a stoichiometric (111) spinel surface as $(1/2)16d-1(\text{Li}_{1/6}\text{Ti}_{5/6})/8a(\text{Li})/32e(\text{O})/\dots$, as discussed for LiMn_2O_4 [31]. Compared with this configuration, the type (A) and type (B) surfaces are apparently oxygen-deficient (cation-rich) and oxygen-rich (cation-deficient), respectively. We can also consider the other surface structure with termination of the 16d-1 layer as the type (C) surface, which is also cation-rich. In the previous theoretical studies of spinel (111) surfaces of Fe_3O_4 [25,26] and Co_3O_4 [27], the surface structure models corresponding to the present type (B) and type (C) surfaces are dealt with as oxygen-rich and oxygen-deficient ones, respectively. Density-functional theory (DFT)

calculations indicated the relative stability of the tetrahedral-cation terminated surface, corresponding to the type (B) structure, in Fe_3O_4 [26] and Co_3O_4 [27] in a wide range of oxygen chemical potentials, which is consistent with the recent LEED and STM observations of $\text{Fe}_3\text{O}_4(111)$ and $\text{Co}_3\text{O}_4(111)$ surfaces [24,28]. Note, here, that the present type (B) $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surface and the $\text{Fe}_3\text{O}_4(111)$ and $\text{Co}_3\text{O}_4(111)$ surfaces are cation-terminated and oxygen-rich as the surface region stoichiometry.

The relative stability between the type (A) and type (B) $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surfaces should depend on the atmosphere. It is considered that the type (B) (oxygen-rich) surface is more stable than the type (A) surface in oxygen-rich atmosphere. The partial occurrence of the type (A) surface in STM sample should be caused by the high-temperature treatment in UHV, namely by the oxygen-poor atmosphere (low oxygen chemical potential), which should stabilize the type (A) (oxygen-deficient) surface relatively. In other words, in the ambient atmosphere, the type (A) surface may disappear, which is consistent with only one type of (111) terrace in the AFM sample. Of course, we cannot conclude the presence of the type (B) surface in the AFM sample without detailed examination, while the surface structure of the AFM sample should have some relation with the type (B) surface structure.

However, the MEIS results indicate the relative stability of the type (B) surface even in the oxygen-poor atmosphere. In any cases, the relative stability of the type (B) surface in a wide range of oxygen chemical potentials, and the possible partial stabilization of the type (A) surface after the high-temperature treatment in UHV are consistent with the theoretical and experimental studies of $\text{Fe}_3\text{O}_4(111)$ and $\text{Co}_3\text{O}_4(111)$ surfaces.

There exists a problem of excess electrons or holes in off-stoichiometric oxide surfaces on the other hand. If we consider that all the valence electrons of cation atoms are transferred to neighboring oxygen atoms, leaving each oxygen atom as $2e$ state, there are 2.75 excess electrons per 1×1 period in the type (A) surface, and there is a deficiency of 1.75 electrons per 1×1 period in the type (B) surface. If we assume a simple surface band structure with the valence band mainly consisting of oxygen $2p$ orbitals and the conduction band consisting of cation valence orbitals, the type (A) surface should have 2.75 excess electrons in the conduction band, leaving the surface cation atoms partially metallic (reduced), while the type (B) surface should have 1.75 holes at the valence-band top consisting of surface oxygen orbitals.

We think that the minority portion of the (111) terraces with the type (A) structure is possibly metallic due to the excess electrons so as to attain local charge neutrality. About the presence of holes in the valence band for the majority portion of the (111) terraces with the type (B) structure, however,

we cannot make definitive conclusion. This is because, there are rare examples of oxygen-rich polar surface surfaces in usual oxides [29,32], except for some transition metal oxides such as Fe_3O_4 and Co_3O_4 where the simple views of charge transfer or excess holes do not seem to be applicable, due to complex nature of oxygen-metal p-d hybridization and *d* electrons [25-27]. Thus for the stabilization of the present type (B) surface structure in real $\text{Li}_4\text{Ti}_5\text{O}_{12}$ samples, there may exist some additional effects so as to remove holes, such as adsorption of hydrogen or water molecules [25,26], surface oxygen vacancies or reconstruction, or excess Ti contents in the 16d-2 layer just below the surface oxygen layer to provide larger numbers of valence electrons. In any case, all these issues should be examined in the near future.

Finally, the oxygen-rich $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface structure may play an important role in the surface chemical reactivity. The presence of holes in the surface oxygen band should have remarkable effects on the oxidation of organic molecules in liquid electrolytes in contact with the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surface. It is reported that CO_2 gases are generated from organic liquid electrolytes by reaction with $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surfaces without applying electric power [33]. This may be caused by the redox interactions between holes and organic molecules. Of course, applying the electric power may also affect the formation of or removal of such holes at the oxygen-rich surfaces. In any case, the present knowledge should contribute to the investigation of atomistic mechanisms of various phenomena at the $\text{Li}_4\text{Ti}_5\text{O}_{12}$ /electrolyte interfaces.

4. Conclusions

We investigated the atomistic structure of a $\text{Li}_4\text{Ti}_5\text{O}_{12}$ (111) surface by means of STM, AFM and MEIS analyses. In the STM, there are two kinds of (111) surface terraces with apparent hexagonal atomic arrangements with interatomic distance of 0.6 nm, bounded by steps with two different step heights. We propose that the major portion of the terraces is an oxygen-rich surface, consisting of the sequence as $8a(\text{Li})/32e(\text{O})/16d-2(\text{Li}_{1/6}\text{Ti}_{5/6})/32e(\text{O})$ and that the minor portion of the terraces is a cation-rich surface, consisting of the sequence of $8a(\text{Li})/16d-1(\text{Li}_{1/6}\text{Ti}_{5/6})/8a(\text{Li})/32e(\text{O})$. The present view was clearly supported by the MEIS analysis to provide reliable information on the T-atom distribution in the surface and sub-surface layers. The major oxygen-rich $\text{Li}_4\text{Ti}_5\text{O}_{12}$ surface should have electronic holes, which is a possible origin of the CO_2 -gas generation via redox interaction with electrolyte molecules.

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